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http://hdl.handle.net/2268/30144

http://orbi.ulg.ac.be/bitstream/2268/30144/1/EII_VLTI_School_J_Surdej.pdf

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"In terms of angular resolution, two small apertures distant of B are equivalent to a single large aperture of diameter B"



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$I(\zeta,\eta) = \iint PSF(\zeta - \zeta',\eta - \eta') O(\zeta',\eta') d\zeta' d\eta' = PSF(\zeta,\eta) \otimes O(\zeta,\eta)$

$TF(I(\zeta,\eta))(u,v) = TF(PSF(\zeta,\eta))(u,v) TF(O(\zeta,\eta))(u,v)$

$\mathbf{B}_{\mathbf{u}} = \mathbf{u} \cdot \lambda, \mathbf{B}_{\mathbf{v}} = \mathbf{v} \lambda$

$TF(PSF(\zeta,\eta))(u,v) = \iint A^*(x,y) A (x+u,y+v) dx dy$

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- Introduction
- 2 Reminders
- 3 Brief history of stellar diameter measurements
- 4 Interferometry with two independent telescopes
- 5 Light coherence (Zernicke-van Cittert theorem)
- 6 Examples of optical interferometers
- 7 Results
- 8 Three important theorems (Fundamental, Convolution and Wiener-Khintchin theorems)!

Introduction



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- 2 Reminders
- 2.1. Representation of an electromagnetic wave



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- 2.1. Representation of an electromagnetic wave
- $E = Re\{a exp[i2\pi(vt z / \lambda)]\}$ (2.1.3)
- E = Re{ a exp[-i ϕ] exp[i2 π vt]} (2.1.4) where $\phi = 2\pi z / \lambda$. (2.1.5)
- $E = a \exp[-i \phi] \exp[i2\pi vt]$

(2.1.6)

An introduction to optical/IR interferometry 2.1. Representation of an electromagnetic wave $E = A \exp[i2\pi vt]$ (2.1.7)with $A = a exp[-i \phi]$ (2.1.8)

 $v \sim 6 \ 10^{14} \, \text{Hz}$ for $\lambda = 5000 \, \text{\AA}$

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2.1. Representation of an electromagnetic wave

$$\langle E^2 \rangle = \lim_{T \to \infty} \frac{1}{2T} \int_{-T}^{+T} E^2 dt$$

(2.1.9)

$$\langle E^2 \rangle = a^2 / 2 \tag{2.1.10}$$

 $I = A A^* = |A|^2 = a^2$.

(2.1.11)

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2.2. The Huygens-Fresnel principle



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2.2. The Huygens-Fresnel principle







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2.2. The Huygens-Fresnel principle



Brief history of stellar diameter measurements



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Brief history of stellar diameter measurements

b) Newton:

$$V_{\odot} - V = -5 \log (z / z_{\odot}), \quad (3.1)$$
$$\Delta = 2 R_{\odot} / z , \quad (3.2)$$
$$\Delta \sim 2 10^{-3}, \quad (8 10^{-3}), \quad (3.3)$$

c) Fizeau-type interferometry

4 Interferometry with two independent telescopes



4 Interferometry with two independent telescopes

b) Fizeau ... the father of stellar interferometry (1868) If $\Delta \ge \phi/2 = \lambda / (2B)$, (4.7) fringe disappearance!

Fringe visibility:

$$\upsilon = \left(\frac{I_{\max} - I_{\min}}{I_{\max} + I_{\min}}\right)$$



4 Interferometry with two independent telescopes

b) Fizeau ... the father of stellar interferometry (1868)



- 4 Interferometry with two independent telescopes
- b) Fizeau ... the father of stellar interferometry (1868)



Stéphan, 1873 Δ << 0,16''

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Marseille 80 cm telescope



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4 Interferometry with two independent telescopes
 b) Fizeau ... the father of stellar interferometry (1868)



• Michelson and Pease (1920)



4 Interferometry with two independent telescopes
 b) Fizeau ... the father of stellar interferometry (1868)

- Anderson
- Brown and Twiss (1956)
- Radio Interferometry (1950)



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- 5 Light coherence
- 5.1 Quasi monochromatic light (waves and wave groups)



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- 5 Light coherence
- 5.1 Quasi monochromatic light (waves and wave groups)

$$V(z,t) = \int_{v-\Delta v}^{v+\Delta v} \frac{a(v') \exp(i2\Pi(v't-z/\lambda'))dv'}{(5.1.1)} (5.1.1)$$

$$I = \langle V(t)V^{*}(t) \rangle \qquad \frac{\exp(i2\Pi(vt-z/\lambda))}{(5.1.2)} \exp(i2\Pi(vt-z/\lambda)) (5.1.2)$$

$$V(z,t) = A(z,t) \exp(i2\Pi(vt-z/\lambda)) (5.1.3)$$

$$A(z,t) = \int_{v-\Delta v}^{v+\Delta v} a(v') \exp(i2\Pi((v'-v)t-z(1/\lambda'-1/\lambda)))dv' (5.1.4) (5.1.4)$$

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- 5 Light coherence
- 5.1 Quasi monochromatic light (waves and wave groups)



- 5 Light coherence
- 5.1 Quasi monochromatic light (waves and wave groups)

 $λ_0 = 2.2 μm$ λ ∈ [2.07 ; 2.33]μmΔ λ = 0.13 μm



- 5 Light coherence
- 5.2 Fringe visibility



- 5 Light coherence
- 5.2 Fringe visibility

$$I_{q} = I + I + 2 I \operatorname{Re}\{\gamma_{12}(\tau)\}$$
(5.2.5)

$$\gamma_{12}(\tau) = \langle V_1^*(t) V_2(t-\tau) \rangle / I$$
(5.2.6)

$$\gamma_{12}(\tau) = \langle A_1^*(z,t) A_2(z,t-\tau) \rangle \exp(-i2\Pi \nu \tau) / I$$
(5.2.7)



$$\gamma_{12}(\tau) = |\gamma_{12}(\tau=0)| \exp(i\beta_{12} - i2\Pi \nu \tau)$$

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(5.2.8) 33

- 5 Light coherence
- 5.2 Fringe visibility

$$I_{q} = I + I + 2I |\gamma_{12}(0)| \cos(\beta_{12} - 2\Pi \nu \tau)$$
(5.2.9)

$$\boldsymbol{\upsilon} = \left(\frac{\boldsymbol{I}_{\max} - \boldsymbol{I}_{\min}}{\boldsymbol{I}_{\max} + \boldsymbol{I}_{\min}}\right) = \left|\boldsymbol{\gamma}_{12}(0)\right|$$

(5.2.10)

- 5 Light coherence
- 5.3 Spatial light coherence

??
$$\gamma_{12}(\tau=0) = \langle V_1^*(t) V_2(t) \rangle / I_{(5.3.1)}$$

$$\begin{cases} V_{1}(t) = \sum_{i=1}^{N} V_{i1}(t) \\ V_{2}(t) = \sum_{i=1}^{N} V_{i2}(t) \\ I_{q}? \end{cases}$$

$$S = \Sigma dS_{i}$$
for i = 1, N
$$P_{1}$$

$$V_{1}(t)$$

$$I_{q}?$$

$$I_{q}?$$

$$I_{q}?$$

$$P_{2}$$

$$V_{2}(t)$$

$$V_{2}(t)$$

$$V_{12}(0) = \left[\sum_{i=1}^{N} \langle V_{2}^{*}V_{i} \rangle + \sum_{j=1}^{N} \langle V_{j}^{*}V_{j} \rangle + \sum_{j=1}^{N}$$

5 Light coherence

5.3 Spatial light coherence

$$\gamma_{12}(0) = \left[\sum_{i=1}^{N} \langle V_{i1}^* V_{i2} \rangle + \sum_{i\neq j}^{N} \langle V_{i1}^* V_{j2} \rangle\right] / I_{(5.3.3)}$$

$$V_{i1}(t) = \left(\frac{a_i(t - \gamma_{i1}/c)}{\gamma_{i1}}\right) \exp\left\{i2\Pi v(t - \gamma_{i1}/c)\right\}_{(5.3.4)}$$
$$V_{i2}(t) = \left(\frac{a_i(t - \gamma_{i2}/c)}{\gamma_{i2}}\right) \exp\left\{i2\Pi v(t - \gamma_{i2}/c)\right\}$$

$$V_{i1}^{*}(t)V_{i2}(t) = \left| a_{i}(t - r_{i1}/c) \right|^{2} / (r_{i1}r_{i2}) \exp\{-i2\Pi v(r_{i2} - r_{i1})/c\}$$
(5.3.5)
as long as:
$$\left| r_{i1} - r_{i2} \right| \le c/\Delta v = \lambda^{2}/\Delta \lambda = \ell$$
(5.3.6)

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- 5 Light coherence
- 5.3 Spatial light coherence

$$I(s)ds = \left|a_i(t-r/c)\right|^2$$

$$\gamma_{12}(0) = \int_{S} \frac{I(s)}{r_{1}r_{2}} \exp\{-i2\Pi(r_{2}-r_{1})/\lambda\} ds / I$$
(5.3.8)

!!! Theorem of Zernicke-van Cittert !!!

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(5.3.7)

- 5 Light coherence
- 5.3 Spatial light coherence



 $|\mathbf{r}_2 - \mathbf{r}_1| = |\mathbf{P}_2 \mathbf{P}_i - \mathbf{P}_1 \mathbf{P}_i| = |-(\mathbf{X}^2 + \mathbf{Y}^2) / 2 \mathbf{Z}' + (\mathbf{X} \boldsymbol{\zeta} + \mathbf{Y} \boldsymbol{\eta})|$ (5.3.9)

where $\zeta = X' / Z'$ and $\eta = Y' / Z'$

(5.3.10)

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- 5 Light coherence
- 5.3 Spatial light coherence

$$\gamma_{12}(0, X/\lambda, Y/\lambda) = \exp(-i\phi_{X,Y}) \frac{\int I(\zeta, \eta) \exp\{-i2\Pi(X\zeta + Y\eta)/\lambda\} d\zeta d\eta}{\iint I(\zeta, \eta') d\zeta' d\eta'}$$
(5.3.11)

$$I'(\zeta,\eta) = I(\zeta,\eta) / \iint_{S} I(\zeta',\eta') d\zeta' d\eta'$$
(5.3.12)

Setting $u = X/\lambda$, $v = Y/\lambda$:

$$\gamma_{12}(0,u,v) = \exp(-i\phi_{u,v}) \iint_{S} I'(\zeta,\eta) \exp\{-i2\Pi(u\zeta+v\eta)\} d\zeta d\eta$$
(5.3.13)

$$I'(\zeta,\eta) = \iint \gamma_{12}(0,u,v) \exp\{i2\Pi(\zeta u + \eta v)\}d(u)d(v)$$
(5.3.14)

5.4 Fourier transform (cf. Léna 1996) 5.4.1 Definitions:

$$TF_{-}f(s) = \int_{-\infty}^{\infty} f(x) e^{-2i\pi sx} dx,$$

$$f(x) = \int_{-\infty}^{\infty} TF - f(s) e^{2i\pi sx} ds$$

$$\int_{-\infty}^{\infty} \left| f(x) \right|^2 dx.$$

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(5.4.1)

(5.4.2)

(5.4.3)

5.4 Fourier transform (cf. Léna 1996)
 5.4.1 Definitions: Generalisation:

$$TF _ f(\vec{w}) = \int_{-\infty}^{\infty} f(\vec{r}) e^{-2i\pi \vec{r} \cdot \vec{w}} d\vec{r}$$
 (5.4.4)

5.4.2 Some properties: a) Linearity: $TF_(af) = a TF_f$, $a \in \Re$, a being a constant,

(5.4.5)

$TF_{(f+g)} = TF_{f} + TF_{g}.$

(5.4.6)

5.4 Fourier transform (cf. Léna 1996)
 5.4.2 Some properties: b) Symmetry & parity:



5.4 Fourier transform (cf. Léna 1996)
 c) Similitude:

$$TF_{(f(ax))(s)} = |a|^{-1} TF_{(f(x))(s/a)},$$

(5.4.9)

where $a \in \Re$, is a constant.

d) Translation: TF_(f(x - a))(s) = $e^{-2i\pi as}$ TF_(f(x))(s)

(5.4.10)

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- 5.4 Fourier transform (cf. Léna 1996)
 e) Derivation:
- $$\label{eq:transform} \begin{split} \mathsf{TF}_{df/dx}(s) &= 2i\pi s \; \mathsf{TF}_{f}(s), \; \mathsf{TF}_{d^nf/dx^n}(s) = (2i\pi s)^n \\ & \mathsf{TF}_{f}(s). \end{split}$$
- 5.4.3 Some important cases (one dimension):a) Door function:

 $\Pi(x) = 1 \text{ if } x \in]-1/2, 1/2[,$

(5.4.12)

= 0 if x ∈]-∞, -1/2] or x ∈ $[1/2, \infty]$.

5.4 Fourier transform (cf. Léna 1996)

The door function and its Fourier transform (cardinal sine)



 $TF_{(\Pi(x))(s)} = sinc(s) = sin(\pi s) / \pi s.$

(5.4.13)

 $TF_{(\pi(x/a))(s)} = |a| \sin(as) = |a| \sin(\pi as) / \pi as.$ (5.4.14)

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5.4 Fourier transform (cf. Léna 1996)
 b) Dirac distribution:

$$\delta(x) = \int_{-\infty}^{\infty} e^{2i\pi sx} ds$$

(5.4.15)

its Fourier transform is thus unity (= 1) in the interval]- ∞ , ∞ [.

- 5 Light coherence
- 5.5 Aperture synthesis

$$\upsilon = \left| \gamma_{12}(0, B/\lambda) \right| = \left| \frac{\sin(\Pi Bb/\lambda z')}{\Pi Bb/\lambda z'} \right|$$
(5.5.2)

 $\Pi Bb / \lambda z' = \Pi (5.5.3)$

 $\Delta \sim \lambda / B$, for a (5.5.4) rectangular source.

 $\Delta \sim 1.22 \lambda / B$, for^(5.5.5) a circular source !



- 5 Light coherence
- 5.5 Aperture synthesis

Exercises (...): point-like source?, double pointlike source with a flux ratio = 1?, gaussian-like source?, uniform disk source?, ...

Case of a double point-like source with a flux ratio = 1



Case of a double point-like source with a flux ratio 0.7/0.3



Variation of the fringe contrast as a function of the angular separation between the two stars:



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If the source is characterized by a uniform disk light distribution, the corresponding visibility function is given by

$$\upsilon = \left(\frac{I_{\text{max}} - I_{\text{min}}}{I_{\text{max}} + I_{\text{min}}}\right) = \left|\gamma_{12}(0)\right| = TF(I) = \frac{2J_1(\pi\theta_{UD}B/\lambda)}{\pi\theta_{UD}B/\lambda}$$









For the case of the Sun:

 $d(\mu) = 7.2 \text{ or } 14.4 \ \mu \Rightarrow \sigma = 2.44 \ \lambda / d = 7.8^{\circ} \text{ or } 3.9^{\circ}$

See the masks!



First fringes for the Sun: 9/4/2010

 $B = 29.4 \mu$ d = 11.8 μ







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OVLA_Sun_2











Interferometric observations on 10/4/2010 of Procyon, Mars and Saturn, using the 80cm telescope at Haute-Provence Observatory and adequate masks ...







Procyon B = 12 mm d = 2 mm



Mars B = 12 mmd = 2 mm



Saturn B = 4 mmd = 2 mm



Saturn B = 12 mmd = 2 mm




6 Some examples of optical interferometers





First fringes with I2T



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6 Some examples of optical interferometers



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6 Some examples of optical interferometers





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http://www.aeos.ulg.ac.be/HARI/

7 Some results

Star	Spectral type	Luminosity class	Angular diameter × 10 ⁻³ seconds of arc 20		
a Boo	K2	Giant			
α Tau	K5	Giant	20		
a Sco	M1-M2	Super-giant	40		
B Peg	M2	Giant	21		
o Cet	M6e	Giant	47		
α Ori	M1-M2	Super-giant variable	34→47		
Table 2.1.	Stars measure Free	ared with Michel om Pease (1931).	son's interferometer.		

7 Some results

NOM SPECTRE	DIAMÈTRE	MESURÉ		TEMPÉRATURE EFFECTIVE		DISTANCE	
	SPECTRE	λ 0,55 μm en ms. d'arc	λ = 2,2 μm en ms. d'erc	₽ √R ⊚	λ = 0.55 μm en degrés Kelvin	λ - 2,2 μm en degrés Kelvin	en persecs (1 pc - 3,25 el
a Cas	KOII	54±0.6		26±8	4700 ± 300	01FR - 22	45±9
§ And	MOIII	13.2 ± 1.7	14.4 ± 0.5	33±9	3900 ± 260	3711±64	23±3
y And	10311	6.8±0.5		50 ± 14	4650 ± 250	10155610255101531	75±15
a Per	F51b	29194		\$5 主9	7000±500	1 1	176±6
a Cya	A2la	27±03		145±45	\$200 ± 600	1	500±100
a Ari	K260	7,4±1		15±5	4300 ± 350		23 ± 4
ß Gem	KONI	7.8±0.8		#±2	4900 ± 220		11±1
β Umì	×411	8.9±1		30±9	4229±300		31±11
y Dra	KSHI	8.7±0.8	10.2±1.4	45±10	4300 ± 230	3960±270	59±21
5 Dra	Gani	3.8±0.3	i	15 ± 5	4530 ± 220	122772481 R50	36 ± 8
µ Gem	MON		14.6 2 0.8	94±30		3960±95	60±15
a Tau	KSHI		20.7±0.4	47:#7		3904±34	21±3
a Boo	K2111		21.5±1.2	25±4	000000000000000	4240±120	11±2
a Aur.	GSIII	8.0±1.2	100000000000000000000000000000000000000	11.7:22	6400±200		13.7±0.6
a Auro	GOIII	4.8 ± 1.5		7.1±2	5950 ± 200	i i	13.7±0.6
a Lyr	AQV	30±02		2.6±0.2			8.1±0.3

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....

 6 Some examples of optical interferometers Interferometry to-day is:

Very Large Telescope Interferometer (VLTI)

- 4 x 8.2m UTs
- 4 x 1.8m ATs
- Max. Base: 200m







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6 Some examples of optical interferometers



 6 Some examples of optical interferometers Interferometry to-day is also:

The CHARA interferometer

6 x 1m
telescopes
Max. Base:
330m





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 6 Some examples of optical interferometers Interferometry to-day is also:

Palomar Testbed Interferometer (PTI)

3 x 40cm
telescopes
Max. Base:
110m



6 Some examples of optical interferometers

Interferometry to-day is also:

Keck interferometer

• 2 x 10m telescopes

• Base: 85m



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Closure phases – what are these?

• Measure visibility phase (Φ) on three baselines simultaneously.

• Each is perturbed from the true phase (ϕ) in a particular manner: $\Phi_{12} = \phi_{12} + \varepsilon_1 - \varepsilon_2$ $\Phi_{23} = \phi_{23} + \varepsilon_2 - \varepsilon_3$ $\Phi_{31} = \phi_{31} + \varepsilon_3 - \varepsilon_1$



• Construct the linear combination of these: $\Phi_{12} + \Phi_{23} + \Phi_{31} = \phi_{12} + \phi_{23} + \phi_{31}$ The error terms are antenna dependent – they vanish in the sum. The source information is baseline dependent – it remains. We still have to figure out how to use it!

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Closure phase is a peculiar linear combination of the true Fourier phases:

 In fact, it is the argument of the product of the visibilities on the baselines in question, hence the name triple product (or bispectrum):

 $V_{12}V_{23}V_{31} = |V_{12}| |V_{23}| |V_{31}| \exp(i2\pi [\Phi_{12} + \Phi_{23} + \Phi_{31}]) = T_{123}$ - So we have to use the closure phases as additional constraints In some nonlinear iterative inversion scheme.

Star

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 6 Some examples of optical interferometers Interferometry to-day is also:

Nullin interferometry

- Measurement of « stellar leakage »
- Allow to resolve stars with a a small size interferometer



 6 Some examples of optical interferometers Interferometry to-day is also:



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6 Other examples of interferometers: ALMA



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6 Other examples of interferometers: DARWIN



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8 Three important theorems ... and some applications

8.1 The fundamental theorem

8.2 The convolution theorem

8.3 The Wiener-Khintchin theorem

Réf.: P. Léna; Astrophysique: méthodes physiques de l'observation (Savoirs Actuels / CNRS Editions)

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8.1 The fundamental theorem

 $a(p,q) = \mathsf{TF}_{A(x,y)}(p,q),$ $a(p,q) = \int_{R^2} A(x,y) \exp[-i2\pi(px+qy)] dx dy,$

2010

with

 $p = x'/(\lambda f)$ $q = y'/(\lambda f)$



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8.1 The fundamental theorem

The distribution of the complex amplitude a(p,q) in the focal plane is given by the Fourier transform of the distribution of the complex amplitude A(x,y) in the entrance pupil plane.

8.1 The fundamental theorem Application: Point Spread Function determination



$$A(x,y) = A_0 P_0(x,y),$$
 (8.1.1)

$$P_0(x,y) = \Pi(x / a) \Pi(y/a).$$
 (8.1.2)

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8.1 The fundamental theorem $a(p,q) = TF [A(x,y)](p,q) = \int_{-a/2}^{a/2} \int_{-a/2}^{a/2} A_0 \exp[-i2\pi(px+qy)]dxdy \quad (8.1.3)$

$$a(p,q) = A_0 \int_{-a/2}^{a/2} \exp[-i2\pi px] dx \int_{-a/2}^{a/2} \exp[-i2\pi qy] dy$$
(8.1.4)

 $a(p,q) = A_0 a^2 [sin(\pi pa) / (\pi pa)] [sin(\pi qa) / (\pi qa)].$ (8.1.5)

 $\frac{i(p,q) = a(p,q) a^{*}(p,q) = |a(p,q)|^{2} = |h(p,q)|^{2} =$ $= i_{0} a^{4} [sin(\pi pa) / (\pi pa)]^{2} [sin(\pi qa) / (\pi qa)]^{2}.$ (8.1.6)

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8.1 The fundamental theorem Application: Point Spread Function determination



 $\Delta p = \Delta q = \Delta x'/(\lambda f) = \Delta y'/(\lambda f) = 2/a \Rightarrow \Delta \phi_{x'} = \Delta \phi_{y'} = 2\lambda/a \quad (8.1.7)$ 2010 VLTI school - Porquerolles

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BESSEL FUNCTIONS (REMINDER)

Integral representation of the Bessel functions

$$\boldsymbol{J}_{0}(x) = \frac{1}{\pi} \int_{0}^{\pi} \cos[x \sin(\vartheta)] d\vartheta \qquad \boldsymbol{J}_{n}(x) = \frac{1}{\pi} \int_{0}^{\pi} \cos[n\vartheta - x \sin(\vartheta)] d\vartheta$$

Undefined integral

$$\int x' J_0(x') dx' = x J_1(x)$$

Graphs of the $J_0(x)$ and $J_1(x)$ functions



Series development (x ~ 0): $J_0(x) = 1 - x^2/2^2 + x^4/(2^24^2) - x^6/(2^24^26^2) + \dots$ $J_1(x) = x/2 - x^3/(2^24) + x^5/(2^24^26) - x^7/(2^24^26^28) + -1$ $J_n(x) = (2 / (\pi x))^{1/2} \cos(x - n\pi/2 - \pi/4) \dots \text{ and when } x \text{ is large!}$

8.1 The fundamental theorem: 2 telescope interferometer



Two coupled optical telescopes: simplified optical scheme (a). Distribution of the complex amplitude for the case of two circular (b) or square (c) apertures and 104 corresponding impulse response (d).

8.1 The fundamental theorem: 2 telescope interferometer

$$h(p,q) = TF(P(x,y)(p,q) = \int_{R^2} P(x,y) \exp[-i2\pi(px+qy)] dxdy \quad (8.1.10)$$

$$h (p,q) = TF(P_0(x+D/2) + P_0(x-D/2))(p,q) =$$

$$TF(P_0(x+D/2))(p,q) + TF(P_0(x-D/2))(p,q) =$$

$$\exp(i\pi D) TF(P_0(x))(p,q) + \exp(-i\pi D) TF(P_0(x))(p,q) =$$

$$(\exp(i\pi D) + \exp(-i\pi D)) TF(P_0(x))(p,q) =$$

$$2\cos(\pi D) TF(P_0(x))(p,q)$$

For the particular case of two square apertures: $i(p,q) = \left| h(p,q) \right|^{2} = 4 \cos^{2}(\pi pD) d^{4} \left(\frac{\sin(\pi qd)}{\pi qd} \right)^{2} \left(\frac{\sin(\pi pd)}{\pi pd} \right)^{2}$

(8.1.12)

(8.1.11)

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Delay lines at the VLTI


What do the VLTI telescope locations look like?





How are those locations related to the *uv* coverage?



Examples of Fourier plane coverage

Dec -15

Dec -65



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How does the uv plane coverage impact imaging?



8.2 The convolution theorem

$$f(x) * g(x) = (f * g)(x) = \int_{R^n} f(x - t)g(t)dt$$



Convolution product of two 1D rectangle functions. A) f(x), B) g(x), C) g(t) and f(x-t); the dashed area represents the integral of the product of f(x-t) and g(t) for the given x offset, D) f(x)*g(x) = (f*g)(x) represents the previous integral as a function of x.

8.2 The convolution theorem

 $e(p,q) = O(p,q) * |h(p,q)|^2$,



$$e(p,q) = \int_{R^2} O(r,s) \left| h(p-r,q-s) \right|^2 drds$$

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8.2 The convolution theorem

For the case of a point-like source:

 $O(p,q) = E \, \delta(p,q),$

(8.2.1)

 $[\delta(x) = 0 \text{ if } x \neq 0, \delta(x) = \infty \text{ if } x = 0] \text{ and}$ (8.2.2)

 $e(p,q) = O(p,q) * |h(p,q)|^2 = E \delta(p,q) * |h(p,q)|^2 = E |h(p,q)|^2 (8.2.3)$

8.2 The convolution theorem

More generally, since TF_(f * g) = TF_(f) TF_(g).

We find, because $e(p,q) = O(p,q) * |h(p,q)|^2$ (8.2.5)

(8.2.4)

that: $TF_(e(p,q)) = TF_(O(p,q)) TF_(|h(p,q)|^2), (8.2.6)$

and, finally, $O(p,q) = TF^{-1}[TF_{e}(p,q)) / TF_{h}(|h(p,q)|^{2})].$ (8.2.7)

8.2 The convolution theorem

$$O(p,q) = (\lambda^{2} E / \phi^{2}) \Pi (p \lambda / \phi) \Pi (q \lambda / \phi).$$

$$e(p,q) = O(p,q) * |h_{0}(p,q)|^{2}.$$

$$e(p) = O(p) * |h_{0}(p)|^{2},$$

$$e(p) = 2d^{2}(\lambda/\phi)\sqrt{E} \int_{p-\phi/2\lambda}^{p+\phi/2\lambda} \left(\frac{\sin(\pi r d)}{\pi r d}\right)^{2} \cos^{2}(\pi r D) dr$$
(8.2.10)

$$\left(\frac{\sin(\pi r d)}{\pi r d}\right)^{2} \approx \text{Cte sur } [p-\phi/2\lambda, p+\phi/2\lambda], \text{ et } (8.2.11)$$

$$e(p) = 2 \frac{d^{2}(\lambda/\phi)}{\sqrt{E}} \left(\frac{\sin(\pi p d)}{\pi p d}\right)^{2} \int_{p-\phi/2\lambda}^{p+\phi/2\lambda} \cos^{2}(\pi r D) dr.$$

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8.2 The convolution theorem

$$e(p) = 2d^{2} \left(\frac{\sin(\pi p d)}{\pi p d} \right)^{2} \left[O(p) *_{\cos^{2}(\pi p D)} \right], \qquad (8.2.13)$$

$$e(p) = 2d^{2} \left(\frac{\sin(\pi pd)}{\pi pd}\right)^{2} \left[\frac{1}{2} \int_{R} O(p)dp + \frac{1}{2} O(p) * \cos(2\pi pD)\right]$$
(8.2.14)

$$e(p) = A \left[B + \frac{1}{2} \operatorname{Re}(O(p) * \exp(i2\pi pD)) \right],$$
 (8.2.15)

$$A = 2d^{2} \left(\frac{\sin(\pi pd)}{\pi pd}\right)^{2} et \quad B = \frac{1}{2} \int_{R} O(p) dp, \qquad (8.2.16)$$
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8.2 The convolution theorem

$$e(p) = A \left[B + \frac{1}{2} \operatorname{Re}\left(\int_{R} O(r) \exp(i2\pi(p-r)D) dr \right) \right], \qquad (8.2.17)$$

$$e(p) = A \left| B + \frac{1}{2} \cos(2\pi pD) TF_{(O(r))(D)} \right|,$$
 (8.2.18)

$$\gamma(D) = (e_{\text{max}} - e_{\text{min}}) / (e_{\text{max}} + e_{\text{min}}),$$

(8.2.19)

 $\gamma(D) = TF_(O(r))(D) / (2B) = TF_(O(r))(D) / \int O(p) dp.$ (8.2.20)

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8.3 The Wiener-Khintchin theorem The Wiener-Khinchin theorem states that the power spectral density of a wide-sense-stationary random process is the Fourier transform of the corresponding autocorrelation function. In our case, this theorem merely states that the Fourier tranform of the PSF (see Eq. (8.2.7)) is the auto-correlation function of the distribution of the complex amplitude in the pupil plane:

$$\frac{TF(|h(p,q)|^2)}{\int A^*(x,y)} A(x+p,y+q) dx dy$$