Classic disc physics II - simple analytical models

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#### Lectures outline

- Lecture I Basic equations
- Lecture II Steady state solutions Time dependent models Outbursts
- Lecture III Sources of angular momentum transport

• Continuity:

$$\frac{\partial \Sigma}{\partial t} + \frac{1}{R} \frac{\partial}{\partial R} \left( R \Sigma u \right) = 0$$

• In a steady state:

$$\dot{M} = -2\pi R \Sigma u$$

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• Angular momentum conservation:

$$\frac{\partial}{\partial t}(\Sigma\Omega R^2) + \frac{1}{R}\frac{\partial}{\partial R}(\Sigma\Omega R^3 u) = -\frac{1}{R}\frac{\partial}{\partial R}(\nu\Sigma R^3|\Omega'|)$$

• In a steady state:

$$\dot{M}j(R_{
m in})=\dot{M}\Omega R^2-2\pi
u\Sigma R^3\Omega'$$

$$\dot{M} = -2\pi R\Sigma u$$

• Angular momentum conservation:

$$\frac{\partial}{\partial t}(\Sigma\Omega R^2) + \frac{1}{R}\frac{\partial}{\partial R}(\Sigma\Omega R^3 u) = -\frac{1}{R}\frac{\partial}{\partial R}(\nu\Sigma R^3|\Omega'|)$$

• In a steady state:

$$\dot{M}j(R_{
m in})=\dot{M}\Omega R^2+3\pi
u\Sigma R^2\Omega$$

$$\dot{M} = -2\pi R\Sigma u$$

• Angular momentum conservation:

$$\frac{\partial}{\partial t}(\Sigma\Omega R^2) + \frac{1}{R}\frac{\partial}{\partial R}(\Sigma\Omega R^3 u) = -\frac{1}{R}\frac{\partial}{\partial R}(\nu\Sigma R^3|\Omega'|)$$

• In a steady state:

$$3\pi
u\Sigma = \dot{M}\left[1 - \sqrt{rac{R_{
m in}}{R}}
ight]$$

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#### Some comments

• Far from inner boundary

$$\dot{M} = 3\pi \nu \Sigma$$

• If viscosity increases, say, linearly with R, in a steady state  $\Sigma$  decreases linearly

#### Velocity ordering

- Radial velocity:  $u \sim u/R$
- Popular prescription (Shakura & Sunyaev 1973):

$$u = lpha c_{
m s} H = lpha \left( rac{H}{R} 
ight)^2 \Omega R^2$$

• Ordering

$$u = \alpha (H/R)^2 v_{\phi} \ll c_{\rm s} = (H/R) v_{\phi} \ll v_{\phi}$$

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#### Energetics

- During accretion, gravitational binding energy is released
- Dissipated by viscosity
- Rate of viscous dissipation

$$D(R) = \nu \Sigma (R\Omega')^2 = \frac{3}{4\pi} \frac{GM\dot{M}}{R^3} \left[ 1 - \sqrt{\frac{R_{\rm in}}{R}} \right]$$

• If this energy is released as a blackbody by the disc surface

$$\sigma T_{\rm s}^4 = \frac{3}{8\pi} \frac{GM\dot{M}}{R^3} \left[ 1 - \sqrt{\frac{R_{\rm in}}{R}} \right]$$

Note factor 1/2: disc has two sides

#### Simple blackbody disc spectrum

#### • Easy to show:



 $|-\text{SED }\lambda F(\lambda) \sim \lambda^{-(4-2/q)} \sim \lambda^{-(4/3)}$ 



#### More on energetics

• Total energy dissipated by the disc

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$$\int_{R_{\rm in}}^{\infty} 2\pi R D(R) dR = \frac{1}{2} \frac{GM\dot{M}}{R_{\rm in}}$$

– Only half of the potential energy, where did the other half go?

• Energy dissipated at a given radius (width dR)

$$2\pi R D(R) dR = \frac{3}{2} \frac{GM\dot{M}}{R^2} \left[ 1 - \sqrt{\frac{R_{\rm in}}{R}} \right] dR$$

- Another puzzle: only  $GMMdR/2R^2$  comes from gravitational binding energy: where does the rest come from?

#### Time dependent models I - spreading ring



- It is possible to solve analitically in a few simple cases
  - Initially infinitesimally narrow ring (deltafunction) -- Pringle (1981)
  - Constant viscosity

#### Time dependent models I - spreading ring



#### Time dependent models

• Viscous timescale

$$t_{
m visc} \sim R^2 / 
u ~(\sim R/u)$$

- Note: most often this timescale is a strongly increasing function of radius
- Evolution at large distances much slower
- Note: inner parts move in, outer parts move out!
- It is not just accretion, it is spreading!
- Transition from inward to outward motion at



- At large times, most of the disc moves inwards
- An infinitesimally small amount of matter reaches infinitely large distances to carry away the angular momentum

#### Self-similar solutions

- Analytical solutions when v is a power-law (Lynden-Bell & Pringle 1974, Hartmann 1998)
- Here,  $\nu \sim R$
- At small radii, solution approaches steady state, with  $\Sigma \sim R^{-1}$
- At large radii, exponential cut-off

#### Self-similar solutions



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#### Circumstellar discs at very high angular resolution

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#### Timescales

• Orbital motion (dynamical timescale)

$$t_{
m dyn} = \Omega^{-1} = \sqrt{rac{R^3}{GM}}$$

- Vertical hydrostatic balance  $t_{\rm z} = H/c_{\rm s} = \Omega^{-1}$
- Thermal timescale

$$t_{\rm th} = \frac{\Sigma c_{\rm s}^2 / (\gamma(\gamma-1))}{D(R)} = \frac{\Sigma c_{\rm s}^2 / (\gamma(\gamma-1))}{\nu \Sigma (R\Omega')^2} = \frac{4}{9\gamma(\gamma-1)} \frac{1}{\alpha\Omega'}$$

• Ordering

$$t_{\mathrm{dyn}} \sim t_z \ll t_{\mathrm{th}} = (1/\alpha) t_{\mathrm{dyn}} << t_{\mathrm{visc}} = (R/H)^2 t_{\mathrm{th}}$$

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# Thermal/viscous instability and FU Orionis

- FU Orionis: YSOs undergoing sudden outbursts
- In outburst: spectrum well described by active accretion disc (Kenyon & Hartmann)
- In quiescence: consistent with T Tauri

#### Thermal/viscous instability and



From Kenyon & Hartmann (1995)

Circumstellar discs at very high angular resolution

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### Thermal/viscous instability and FU Orionis

- At any radius: heating/cooling balance determines relation between mass flux and  $\Sigma$
- Realistic opacities: gap at 10<sup>3</sup>-10<sup>4</sup> K (Bell & Lin 1994)
- Negative slope in  $M_{dot}$ - $\Sigma$ : instability!
- Disc can only stay either in
  - the upper branch: high accretion, outburst
  - or in the lower one: low accretion, quiescence

## Thermal/viscous instability and FU Orionis



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#### Limit cycle instability in CVs



Courtesy of Patrick Deegan and Graham Wynn

#### Limit cycle in FU Ori?

- Many models proposed:
  - Clarke et al. (1990), Lodato & Clarke (2004)
  - Bell & Lin (1994), Bell et al. (1995)
- Problems:
  - Hard to match different timescales
  - Triggered/untriggered outbursts
  - Required values of  $\alpha \sim 10^{-3}$ -10<sup>-4</sup> very small