

# Classic disc physics

## I - Basics

Giuseppe Lodato  
*University of Leicester*

# Lectures outline

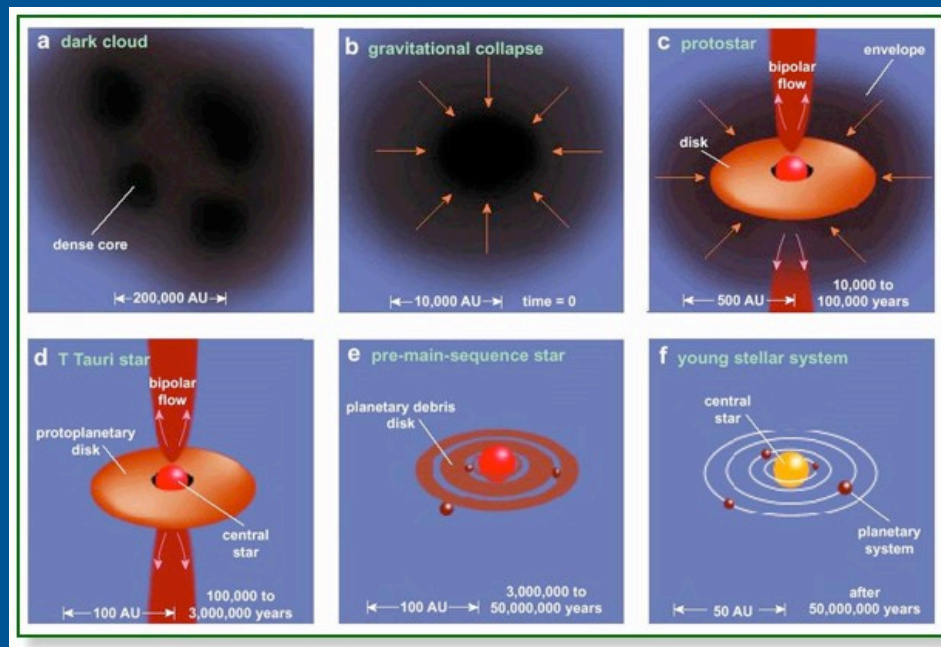
- Lecture I - Basic equations
- Lecture II - Steady state solutions - Time dependent models - Outbursts
- Lecture III - Sources of angular momentum transport

# Discs in protostellar environments

- Protostars are generally surrounded by flattened structures - discs
- General picture of isolated star formation
  - (Things can be significantly more complex in reality! see Bate, Bonnell & Bromm simulations)

# Discs in protostellar environments

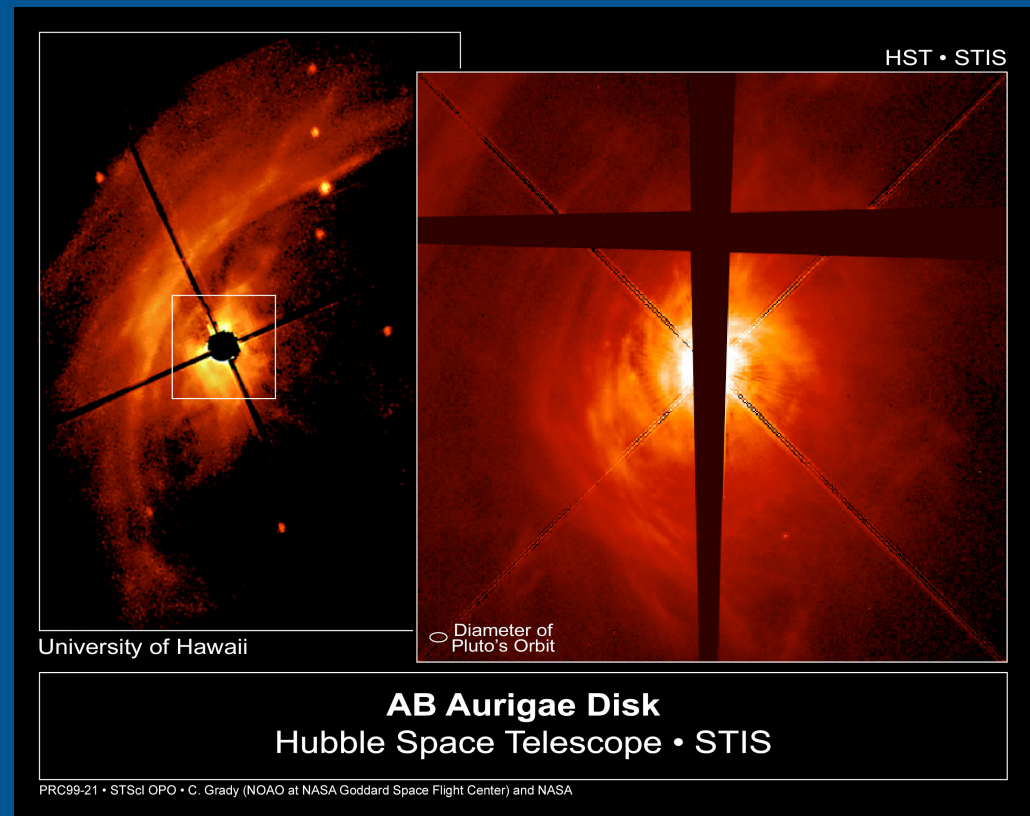
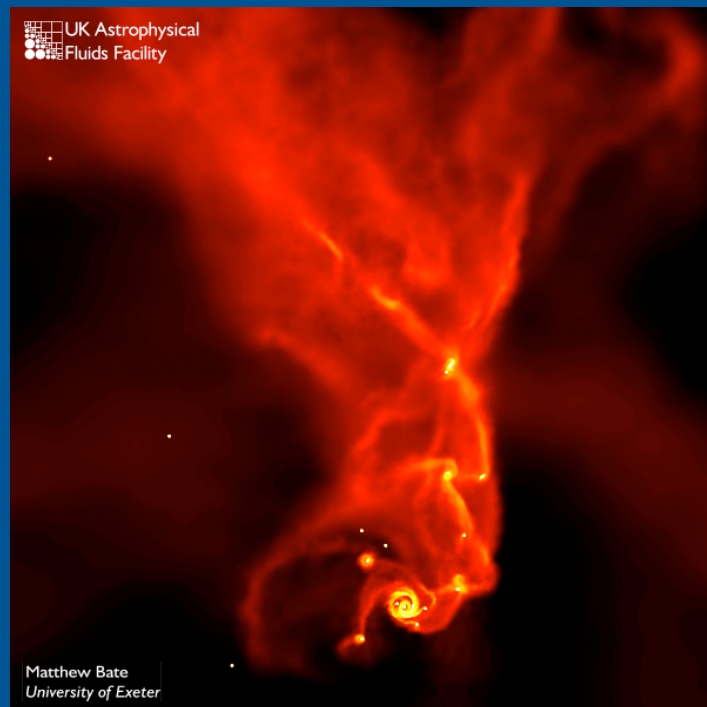
- General picture of isolated star formation



from Blitz

# Discs in protostellar environments

- (Things can be significantly more complex in reality! see Bate, Bonnell & Bromm simulations)



# Discs in protostellar environments

- Discs found at all scales:
  - Around high-mass O-B stars
  - Around intermediate mass star (HAEBE)
  - Around solar mass stars (T Tauri, FU Orionis)
  - Around brown dwarfs
- Protostellar discs are made of:
  - Gas (dominate mass and dynamics)
  - Dust (dominate opacity and thermo-dynamics)

# Main properties: mass

- Obtained from sub-mm observations (where discs are optically thin)
  - **Very uncertain** due to
    - Uncertainties in dust opacity -- grain growth
    - Dust to mass ratio
  - Typically:  $M_{disc} \sim 0.005-0.1 M_*$  (at least in the T Tauri regime)
  - For high-mass stars,  $M_{disc}$  can be a significant fraction of  $M_*$
  - For brown dwarfs?

# Main properties: sizes

- Theory: absolute upper limit: need to reach centrifugal equilibrium

$$t_{\text{dyn}} \sim 2\pi \sqrt{\frac{R^3}{GM}} \sim 10^6 \text{ yrs}$$

$$R \sim 10^4 \text{ AU} (M/M_{\odot})^{1/3}$$

- For BD  $\sim$  a few  $10^3$  AU
- For massive stars  $\sim$  a few  $10^4$  AU
- Observed:
  - High mass stars: pseudo-discs at  $10^5$  AU
  - High mass stars: discs at  $10^4$  AU
  - T Tauri stars: a few 100 AU (big resolved ones)
  - BD: not well determined, maybe comparable to T Tauri



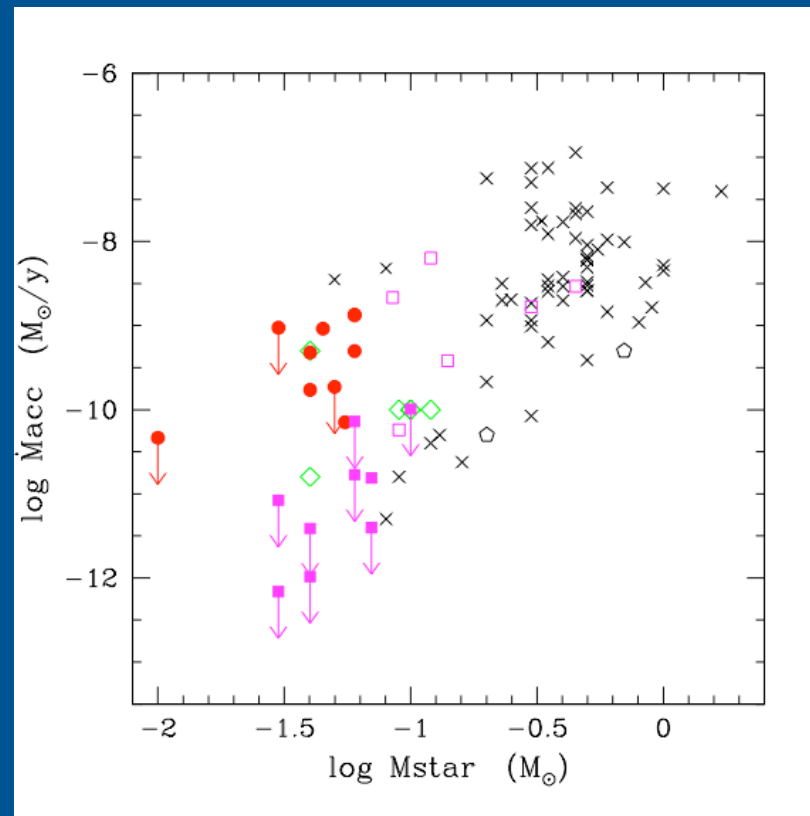
# Main properties

- Mass accretion rate: obtained from a number of methods:
    - strength of emission lines emitted as gas reaches the star
    - veiling of photospheric lines due to accretion shocks
    - Typical numbers:
      - $10^{-5} M_{\text{sun}}/\text{yr}$  for the few high-mass star discs
      - $10^{-7} \text{ -- } 10^{-9} M_{\text{sun}}/\text{yrs}$  for T Tauri stars
      - $10^{-11} M_{\text{sun}}/\text{yrs}$  for brown dwarfs
      - During outbursts:  $10^{-4} M_{\text{sun}}/\text{yrs}$  (FU Orionis)
    - $\dot{M}$  scales as  $M^2$ ?
- (Natta et al. 2004, Clarke & Pringle 2006, Alexander & Armitage 2006)

# Main properties

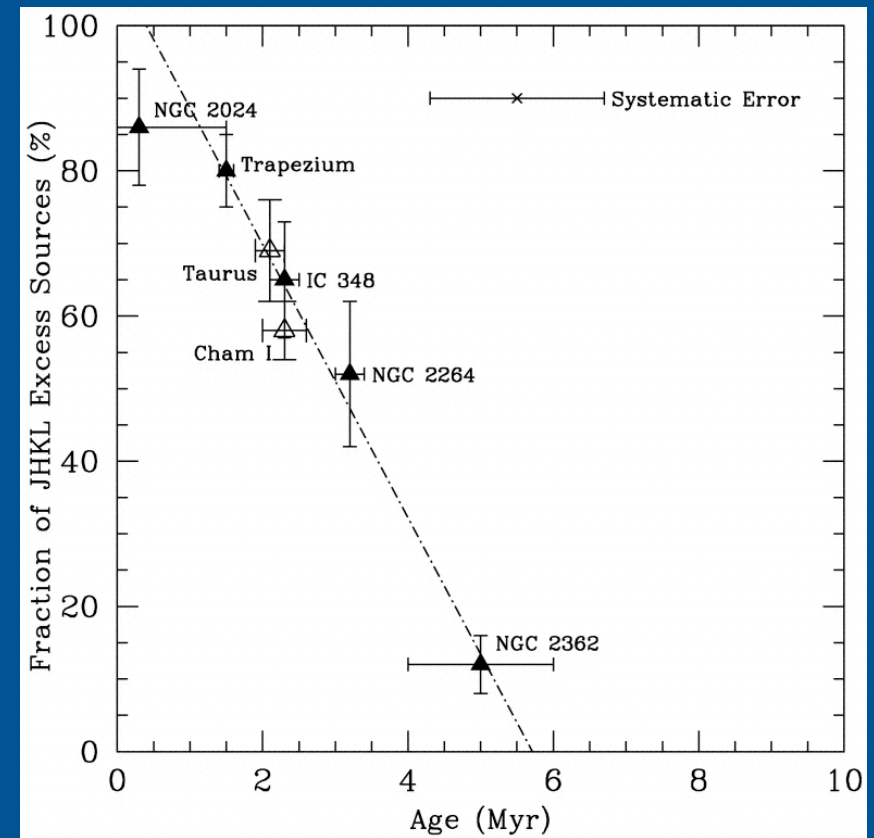
–  $\dot{M}$  scales as  $M^2$ ?

(Natta et al. 2004, Clarke & Pringle 2006, Alexander & Armitage 2006)



# Main properties

- Lifetime:  $10^6$ - $10^7$  years
- Temperature:
  - ranges from  $\sim 10$  K in the outer parts, up to  $\sim 10^3$  K in the hotter inner parts
  - also dependent on state: much higher (up to  $\sim 10^4$  K) during outbursts (e. g. FU Orionis)
  - Temperature determines thickness of the disc (see later)
  - Typically  $H/R \sim 0.1$



Haisch et al. 2001

# Why discs?

- Relatively easy for interstellar gas to get rid of **energy** (e.g., by radiation)
- Much more difficult to remove **angular momentum**
- For a given ang. mom., minimum energy orbit is a circular one
- Reach “circularization radius”, where centrifugal balance holds and orbits become circular
- For a Keplerian potential (assume that accreting gas is much smaller than protostar):

$$R_{\text{circ}} = \frac{L^2}{GM}$$

# Viscous discs

- In a disc, the potential for angular momentum removal is larger
  - Discs are subject to a large number of instabilities that might redistribute angular momentum (Lecture 3)
  - Anisotropies, turbulence generated by instabilities lead to internal stress
  - Initial studies of discs not able to quantify such internal stress and rely on simplified **viscosity** prescription ( $\alpha$ -prescription: Shakura & Sunyaev, 1973)

# A brief history

- The pioneers:
  - Shakura & Sunyaev (1973): viscosity prescription and first steady-state models
  - Lynden-Bell & Pringle (1974): rigorous fluid-dynamical derivation and time-dependent models
- The '80s:
  - Wave propagation in discs:
    - Warps and bending waves (Papaloizou & Pringle 1983)
    - Density waves (Goldreich & Tremaine 1978)
  - Disc-satellite interaction:
    - Lin & Papaloizou (1979)
    - Goldreich & Tremaine (1980)
  - Limit-cycle instability:
    - Clarke, Lin & Papaloizou (1989), Bell & Lin (1994)
- The era of numerical simulations (from the '90s to present)
  - Non-linear disc dynamics, disc instabilities, source of transport
  - The magneto-rotational instability (**MRI**) (Balbus & Hawley 1992)
  - Gravitational instability (**GI**) (Laughlin & Bodenheimer 1994)

# Thin discs

- Main disc property: **being thin!**
- At any radius  $R$ , thickness  $H \ll R$  -- fundamental scaling of all quantities!
- Deal with “surface” quantities: surface density profile  $\Sigma$ ...
- ... or with vertically averaged quantities (for example viscosity  $\nu$ )
- Use cylindrical coordinates

# Basic equations

- Continuity equation:

$$\frac{\partial \Sigma}{\partial t} + \frac{1}{R} \frac{\partial}{\partial R} (R \Sigma u) = 0$$

- Momentum equation:
  - There are actually three components
  - 1) Radial: centrifugal balance
  - 2) Vertical: hydrostatic balance
  - 3) Azimuthal: angular momentum conservation



# Basic equations

- Continuity equation:

$$\frac{\partial \Sigma}{\partial t} + \frac{1}{R} \frac{\partial}{\partial R} (R \Sigma u) = 0$$

- Momentum equation:
  - There are actually three components
  - 1) **Radial: centrifugal balance**

$$v_{\phi}^2 = \Omega^2 R^2 = \frac{GM}{R}$$

# Note

- In order to obtain Keplerian rotation we have neglected several terms in Navier-Stokes eq:
  - pressure terms  $\sim c_s^2/R$
  - Lagrangian derivative  $\sim u^2/R$
  - Is this allowed?
- Note that specific angular momentum in Keplerian discs increases as:

$$j = \Omega R^2 = \sqrt{GM R}$$

# Basic equations

- Continuity equation:

$$\frac{\partial \Sigma}{\partial t} + \frac{1}{R} \frac{\partial}{\partial R} (R \Sigma u) = 0$$

- Momentum equation:
  - There are actually three components
  - 1) Radial: centrifugal balance
  - 2) **Vertical: hydrostatic balance**

# Vertical hydrostatic balance - Keplerian discs

$$\frac{\partial P}{\partial z} = -\frac{GMz}{R^3}\rho$$

$$\frac{\partial \rho}{\partial z} = -\frac{GMz}{c_s^2 R^3}\rho$$

$$\rho = \rho_0 \exp(-z^2/2H^2)$$

$$\frac{H}{R} = \frac{c_s}{V_\phi}$$

- Introduce sound speed (assume it is constant with  $z$ )
- Easy to solve
- Gaussian profile with thickness  $H$
- Thin disc condition equivalent to supersonic rotation:  $c_s \ll V_{rot}$

# Basic equations

- Continuity equation:

$$\frac{\partial \Sigma}{\partial t} + \frac{1}{R} \frac{\partial}{\partial R} (R \Sigma u) = 0$$

- Momentum equation:
  - There are actually three components
  - 1) Radial: centrifugal balance
  - 2) Vertical: hydrostatic balance
  - 3) **Azimuthal: angular momentum conservation**

# Angular momentum conservation

- Navier-Stokes equation

$$\frac{\partial}{\partial t}(\Sigma\Omega R^2) + \frac{1}{R} \frac{\partial}{\partial R}(\Sigma\Omega R^3 u) = -\frac{1}{R} \frac{\partial}{\partial R}(R^2 T_{R\phi})$$

- Viscous stress tensor  $T_{R\phi}$

$$T_{R\phi} = \Sigma \nu R \left| \frac{d\Omega}{dR} \right|$$

# Angular momentum conservation

- Navier-Stokes equation

$$\frac{\partial}{\partial t}(\Sigma \Omega R^2) + \frac{1}{R} \frac{\partial}{\partial R}(\Sigma \Omega R^3 u) = -\frac{1}{R} \frac{\partial}{\partial R}(\Sigma \nu R^3 |\Omega'|)$$

- Viscous stress tensor  $T_{R\phi}$

$$T_{R\phi} = \Sigma \nu R \left| \frac{d\Omega}{dR} \right|$$

# Master disc evolution equation

- Obtained combining continuity equation and angular momentum conservation
- For Keplerian discs,
- .... and after some algebra, get:

$$\frac{\partial \Sigma}{\partial t} = \frac{3}{R} \frac{\partial}{\partial R} \left[ R^{1/2} \frac{\partial}{\partial R} (\Sigma \nu R^{1/2}) \right]$$



# A little puzzle

- Simple kinetic arguments suggest that in a linear shearing viscous flow, momentum is transported from high to low momentum
- Analogously, one would expect ang. mom. in a disc to flow from high to low
- Ang. mom. flux at  $R$ :

$$\propto j(R - dR) - j(R + dR) \propto -\frac{dj}{dR}$$

- However, viscosity prop. to  $-d\Omega/dR$
- How come?

